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Spectroscopic Properties, Electronic Levels, and Crystal Field Parameters of Pentakis(nitrato)europate(III) Ions

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The structural data of some rare-earth pentanitrates, $X_2Ln(NO_3)_5$, are discussed, and the coordination polyhedra around the metal ion are rationalized by use of different approaches. It occurs that the arrangement of the nitrogen atoms is close to a trigonal bipyramid with D_{3h} symmetry. We have tried to find out whether the spectroscopic properties of $X_2Ln(NO_3)_5$ $(X = Ph_3EtP^+, Ph_4As^+, and Me_4N^+)$ reflect this pseudosymmetry, even if the actual site symmetry of the europium ion is C_2 (X = Ph₃EtP⁺ and Ph₄As⁺). IR, Raman, and reflectance spectra are reported and discussed, along with a detailed study of the emission spectra at 77 K. For (Ph₃EtP)₂Eu(NO₃)₅, the crystal field splittings indeed display some patterns that can be interpreted as reflecting the higher, pseudotrigonal symmetry. To prove this more quantitatively, the emission spectrum has been simulated by different sets of crystal field parameters by using the method of descending symmetry. With a D_{3h} Hamiltonian, a rms standard deviation of 18 cm⁻¹ is obtained; the use of these parameters as starting values for a fit with a $C_{2\nu}$ Hamiltonian leads to a final, satisfactory, standard deviation of 9.6 cm⁻¹. The coordination polyhedron of $(Ph_4As)_2Eu(NO_3)_5$ seems to be more distorted, since a calculation with a D_{3h} Hamiltonian yields a rather poor fit; a set of $C_{2\nu}$ parameters reproduces the spectrum with $\sigma = 12$ cm⁻¹. The sets of parameters for the two compounds are rather similar as expected for complexes with similar structure; however, the extreme sensitivity of the emission spectra to the environment of the Eu(III) ion makes it difficult to recognize the pentakis(nitrato) ion solely on the basis of its emission spectrum.

Introduction

Lanthanide nitrates form double salts of different compositions with both inorganic and organic nitrates. All of these salts contain either a hexakis(nitrato) ion²⁻⁴ or a pentakis-(nitrato) ion, $Ln(NO_3)_5^{2-}$. In the solid state, hexanitrates are only isolated with the lighter lanthanoid ions, Ln = La-Sm, whereas pentanitrates are isolated for the entire rare-earth series, provided a judicious choice of the counterion is made. A considerable amount of structural work has been performed on pentanitrates, and crystallographic data are available for $Ln = La, 5,6, Ce, 7 Sm, 8 Eu, 9 Ho, 10 Er^{11}, Y, 12 and Sc. 12$ In all these species but the last one, the rare-earth ion is surrounded by five bidentate nitrates. For Ln = La and Sm, one or two water molecules are also coordinated to the metal ion, which is therefore 11- or 12-coordinated, whereas the other Ln(III) ions are 10-coordinated. The site symmetry of the metallic ion is low, C_2 or even C_1 , which may be the reason why no detailed spectroscopic study of the pentanitrates has been published so far. However, a close examination of the crystal structures reveals that this low symmetry arises sometimes from distortions of a more symmetrical arrangement. Our aim is to test if this latter can be assessed from an analysis of the crystal field splittings. As part of a study of the coordination properties of rare-earth ions, both in solutions^{13,14} and in co-

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Table I. Elemental Analyses of $X_2 Eu(NO_3)_5$

		X		
	Ph₄As	Ph ₃ EtP	Me ₄ N	
% C theor	46.92	45.99	15.75	
% C found	46.77	45.87	15.62	
% H theor	3.28	3.86	3.96	
% H fo und	3.19	3.86	3.93	
% N theor	5.70	6.70	16.07	
% N fou n d	5.56	6.77	16.00	
% P theor % P found		5.93 5.87		
% Eu theor	1 2.3 7	14.55	24.83	
% Eu found	1 2.3 2	14.64	24.72	

ordination compounds,¹⁵⁻¹⁹ we therefore report the results of a spectroscopic investigation of finely powdered samples of $(Ph_4As)_2Eu(NO_3)_5$, $(Ph_3Et)_2Eu(NO_3)_5$, and $(Me_4N)_2Eu$ - $(NO_3)_5$, along with a discussion of their structural properties. The crystal field splittings measured from low-temperature emission spectra (77 K) are reproduced, for the first two compounds, by different sets of crystal field parameters by using the method of the descending symmetry.

Experimental Section

Syntheses. The pentakis(nitrato) complexes were prepared by dissolving 5 mmol of Ph₄AsCl (purum, Fluka) or of Ph₃EtBr (purum, Fluka) and adding a solution of 5 mmol of AgNO3 in 25 mL of acetonitrile. The resulting solution was stirred 15 min in the dark and filtered. A solution of 2.5 mmol of Eu(NO₃)₃·6H₂O in 15 mL of acetonitrile was added to the previous solution or to 5 mmol of Me₄NNO₃ (puriss, Fluka) in 125 mL of acetonitrile, and the reaction mixture was refluxed for 2-6 h. After partial evaporation of the solvent, a few milliliters of alcohol were added and the complexes crystallized slowly at -15 °C. The microcrystalline, colorless compounds were recrystallized in acetonitrile and dried over phosphorus pentoxide, in vacuo. Satisfactory elemental analyses were obtained

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Table II. Eu-O and Nitrate Vibrations of $[Eu(NO_3)_5]^{2-}$ (cm⁻¹)

	Ph	As	Ph,	EtP	Me	4 N
assgnt ^a	IR	R ^b	IR	R	IR	R
Eu-O	176	1 84	173	190		
Eu-O	211		212	212		
Eu-O	228	232				
Eu–O	238	241				
v.		708		705		
v.,	738	740	738	736	730	738
ν,	744		746		745	749
ν.	816		814	820	821	
ν,	1023	1025	1025	1025	1032	1029
V.2	1030	1034	1032	1030	1042	1039
V.	1311	1314	1313	1316	1305	
ν_1	1452		1460		1460	
ν,	1479		1482	1481		
ν_1	1488	1495	1491	1496	1490	1515
$\nu_{2} + \nu_{2}$	1730		1732		1735	
$v_{3} + v_{2}$	1774		1768		1742	
$v_{3}^{2} + v_{2}^{2}$					1780	

^a The nitrate assignments are given in terms of bidentate ions with C_{2v} symmetry. ^b Raman spectra.

for all the complexes and are listed in Table I.

Spectroscopic Measurements. IR spectra were measured as Nujol mulls or KBr pellets by using a Perkin-Elmer 577 spectrometer or a Bruker IFS-113c interferometer. Raman spectra were recorded on a Spex Ramalog-4 spectrometer (647-nm excitation by means of a Spectra-Physics Kr-laser). Reflectance spectra were recorded on a Perkin-Elmer Hitachi spectrometer, Model 340, equipped with a Perkin-Elmer R10-A integrating sphere. The fluorescence spectra were excited either by a Zeiss 450-W xenon high-pressure lamp coupled with a Zeiss M-20 monochromator (band-pass used 4-10 nm) or by a Coherent CR-599 dye laser pumped by a Coherent CR-8 argon laser. The fluorescent light was analyzed at 90° with a 0.5-m Czerny-Turner monochromator from Spex Industries (Model 1870) equipped with a 1200 lines/mm grating blazed at 26° 45', with a maximum resolution of 0.02 nm. Light intensity was measured by a RCA 31034 photomultiplier with a cooled Ga-As photocathode (dark count 8 ips at 1940 V and -20 °C), coupled to an Ortec photon-counting unit. Low-temperature measurements were performed with an Oxford DN-704 cryostat, and the samples were put into a specially designed rotatable probe to allow for precise alignment. Fluorescent lifetimes were measured by using a flash-lamp-pumped CMX-4 dye laser from Chromatix. The output signal of the photomultiplier was fed to a Bruker BC-104 transient recorder coupled to a DL-4000 signal averager and a DL-450 microprocessor from Datalab.

X-ray powder diffraction diagrams were obtained by using Cu K α radiation, a Guinier camera, and an Ital-Structure generator. All the data reported in this paper are averages of two to six determinations on at least two independent samples, prepared in the drybox under a nitrogen atmosphere containing less than 10 ppm of water.

Results and Discussion

Vibrational Spectra. The vibrational spectra of the three pentanitrates are closely related and exhibit three kinds of vibrations arising from (i) the lattice and Eu-O vibrations, (ii) the nitrato groups, and (iii) the counterions. The latter could be identified by comparison with reference compounds containing the same cations. In our case, however, they do not provide relevant information about the structural properties of the pentakis(nitrato) ion and are therefore not discussed here. The situation is different for the nitrate vibrations (Table II): the presence of six bands in the IR spectra indicate that these anions are coordinated to the central ion. The weakness of the Raman emission from the $\nu_4(\mathbf{B}_1)$ stretching mode, the position of $\nu_1(A_1)$ and $\nu_2(A_1)$ at ca. 1470 and 1030 cm⁻¹, respectively, and the presence of two combination bands separated by more than 20 cm⁻¹ in the 1730–1770-cm⁻¹ region²⁰ are all consistent with the presence of bidentate nitrates.

Table III. Fit between Ideal Polyhedra and Crystallographic Data for $X_2 Ln(NO_3)_5^a$

Ln	counterion	$\begin{array}{c} \text{EDod} \\ (C_{2v}) \end{array}$	$\begin{array}{c} \text{BCDod} \\ (D_2) \end{array}$	$\begin{array}{c} \text{BCSAP} \\ (D_{4d}) \end{array}$	
Ce Eu	Ph₃EtP+ Ph₄As+	0.120 0.197	0.209 0.078	0.198 0.143	
Ho Er	NO ⁺ K ⁺	$0.064 \\ 0.065$	0.147 0.173	0.199 0.197	

^a Listed numbers represent the normalized standard deviation $\epsilon/\langle M-L\rangle$.²⁶

In addition, the fact that several modes are split into two or more components points to inequivalence of the nitrate groups, as expected from the structural data.

Reflectance Spectra ($X = Ph_4As^+$ and Ph_3EtP^+). In addition to the weak 4f-4f transitions, these spectra contain bands from the phenyl-ring transitions at 26 340 and 27 000 cm⁻¹ for Ph_4As^+ and 26 600 and 27 500 cm⁻¹ for Ph_3EtP^+ as well as one charge-transfer band at 29 450 cm⁻¹. Taking 1.22 as the optical electronegativity of the Eu(III) ion, corrected for spin pairing and for spin-orbit coupling,²¹ the resulting optical electronegativity of the nitrato groups turns out to be 2.35, a value slightly larger than for the acetylacetonate ion, 2.2.²² The importance of the charge transfer can be estimated from Boudreau's empirical formula,²³ taking 21 290 cm⁻¹ as the ionization energy of $NO_3^{-,24}$ 48 600 cm⁻¹ as the electron affinity of Eu(III),²⁵ and 2.48 Å as the mean donor-acceptor distance (=Eu-O distance⁹). The ionic character of the $Eu-NO_3^-$ bond amounts then to 0.53, a value comparable to that found in $Eu(THD)_3$, where THD is the bidentate ligand C₄H₉COCH-COC₄H₉^{-.25}

Structural Data. In this section, we only discuss the structure of the pentanitrates in which the Ln(III) ion is 10coordinated. The structure of such species can usually be rationalized in terms of three polyhedra: the bicapped square antiprism (BCSAP, D_{4d} symmetry), the bicapped dodecahedron (BCDod, D_2), and the 4A,6B-expanded dodecahedron (EDod, $C_{2\nu}$).²⁶ One way to find which polyhedron fits the structure is to compare the angles and the edges with those of the experimentally determined polyhedron. According to this method,²⁷ the coordination polyhedron for (Ph₄As)Eu-(NO₃)₅ is best described by a BCSAP.⁹ However, in the case of bidentate ligands with small "bites" such as the nitrates, the effective distortion of the experimental polyhedron may be overlooked since the method compares averages of edges and angles that should be equivalent. We have therefore calculated the fit between the X-ray structures and these three ideal polyhedra using Dollase's method, as modified by Drew²⁸ (Table III). According to these data, the coordination around the metal ion in the pentanitrates with Ln = Ce, Ho, and Er can be described in terms of an expanded dodecahedron, whereas $(Ph_4As)_2Eu(NO_3)_5$ seems to have a more symmetrical coordination polyhedron. A third way of describing the coordination around the metal ion is to refer to Bergman and Cotton's postulate.²⁹ If the nitrate ions are considered as occupying one coordination site, the resulting polyhedron should be close to the ideal polyhedron for CN = 10/2 = 5, that is a trigonal bipyramid with D_{3h} symmetry. To compare

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Table IV. Comparison between the Positions of the Nitrogen Atoms of $X_2Ln(NO_3)_5$ and the Trigonal Bipyramid^a

	trig bipyr	Ce'	Ho10	Er11	Eu
α_1 , deg	120	119.3	118.5	119.4	117.4
α_{2} , deg	120	121.5	124.6	120.8	124.3
a, deg	120	119.3	116.7	119.7	117.8
β, deg	180	174.6	178.8	179.1	178.1
γ , deg	9 0	90.0	88.8	88.7	90.0
d(plane), Å		0.00	0.02	0.03	0.00
$R(Ln-N), A^{b}$		2.972 2.94-3.00	2.870 2.77 - 2.90	2.856 2.85-2.89	2.896 2.88–2.91

^{*a*} Atoms N₁, N₂, and N₃ are taken as equatorial ligands; their plane should contain the Ln(III) ion. The three angles in this plane are labeled α_i , ^{*b*} Mean distance and minimum and maximum bond lengths.



Figure 1. X-ray powder diagrams of $(Ph_3EtP)_2Ln(NO_3)_5$: (I) Ln = Ce; (II) Ln = Eu.



Figure 2. Effective ionic radii of 10-coordinated Ln(III) ions in pentanitrates vs. atomic numbers.

the experimental results with the latter, we have considered the geometrical figure described by the five nitrogen atoms of the nitrate groups (Table IV). The deviations with respect to the ideal polyhedron are particularly small; even the Ln-N distances lie in a narrow range, except for Ln = Ho. As expected for complexes with which the chemical bonds are predominantly ionic, Bergman and Cotton's description is therefore quite precise in the case of the lanthanide pentanitrates. In the next section, we show that this higher D_{3h} symmetry is reflected, for (Ph₃EtP)Eu(NO₃)₅, in some features of the crystal field splittings.

Since the crystal structure of $(Ph_3EtP)_2Eu(NO_3)_5$ is not known, we have compared its X-ray powder diagram with that of the corresponding cerium compound. They are almost identical (Figure 1), and therefore, the two pentanitrates can be considered as being isomorphous, so that the Eu(III) ion probably also lies on a site with C_2 symmetry.⁷

Finally, the effective ionic radii of the 10-coordinate Ln(III) ions can be evaluated from the structural data. They are in good agreement with the values obtained from interpolation of ionic radius vs. coordination number plots:³⁰ r = 1.250 (extrapolated value 1.255), 1.178 (1.180), 1.137 (1.138), 1.120 (1.126), and 1.116 (1.117) Å for Ln = Ce, Eu, Y, Ho, and Er, respectively. These ionic radii seem to decrease linearly with the atomic number Z (Figure 2).

Emission Spectrum of $(Ph_3EtP)_2Eu(NO_3)_5$. The excitation and emission spectra were recorded at 77 K. The excitation



Figure 3. Excitation spectrum of $(Ph_3EtP)_2Eu(NO_3)_5$ at 77 K (analyzing wavelength 618 nm).



Figure 4. Emission spectrum of $(Ph_3EtP)_2Eu(NO_3)_5$ at 77 K (excitation wavelength 397 nm (⁵L₆ level)).



Figure 5. Crystal field splittings of $(Ph_3EtP)_2Eu(NO_3)_5$, in cm⁻¹, from excitation and emission spectra at 77 K. The numbers between parentheses are the barycenters of the levels calculated from the observed transitions.

spectrum is presented in Figure 3: there is only one sharp, but extremely weak, ${}^{5}D_{0} \leftarrow {}^{7}F_{0}$ transition (full width at half-height, fwhh = 2 Å), consistent with the presence of one crystallographic site for the Eu(III) ion. Selective excitation of this transition with use of a dye laser only produced a Raman spectrum. The emission spectrum is shown on Figure 4; it is dominated by the hypersensitive transition ${}^{5}D_{0} \rightarrow {}^{7}F_{2}$. The integrated and corrected relative intensities of the transitions from the ${}^{5}D_{0}$ level to the ${}^{7}F_{J}$ manifold are <0.1, 10.3, 100.0, 2.0, 10.4, and 1.6% for J = 0-5, respectively.

All the observed transitions in both the excitation and the emission spectra are listed in Tables S1 and S2 (supplementary material), along with their assignment. A scheme of the crystal field splittings for the 7F_J manifolds is presented in Figure 5. The splitting of the J = 1 and J = 2 level is maximum, which implies a low symmetry. However, a closer look at the crystal field splitting for J = 1 suggests that the sublevels appear as one singlet and one doublet. Moreover, only one main band appears in the emission spectrum for the ${}^5D_0 \rightarrow {}^7F_2$ transition.

Table V. Calculated Crystal Field Parameters (CFP, cm⁻¹) for $(Ph_3EtP)_2Eu(NO_3)_s$ in D_{3h} and C_{2v} Symmetries

		equiv Cl C ₂ :	FP along axis	C_{2v}	sym
CFP	D _{3h} sym	sign +	sign – ^a	sign +	sign – ^a
$\overline{B^2}_0$	-276 ± 16	1	39.4	230 ± 13	192 ± 9
B^2_2		1	70.7	173 ± 9	193 ± 7
B ⁴ 0	-496 ± 24	-1	86.0	-102 ± 24	-206 ± 23
B_{2}^{4}		-1	96.0	-246 ± 19	-441 ± 13
B^{4}_{A}		-2	59.4	259 ± 15	-102 ± 11
B^6	987 ± 27	-2211	1587	-2190 ± 20	1514 ± 18
B ⁶ 2		1075	-1704	1072 ± 15	-1739 ± 9
B ⁶ _		-855	161	-728 ± 15	24 ± 13
B6	$\pm 1996 \pm 16$	-407	-521	-636 ± 14	-544 ± 12
σ	18			11.6	9.7

^a Refers to the B_{6}^{6} sign in D_{3h} symmetry.

Therefore, it seems that the actual site symmetry is a distortion of a more symmetrical arrangement having a C_{3h} or a D_{3h} symmetry. This latter would indeed correspond to the trigonal-bipyramidal arrangement of the nitrogen atoms of the five nitrate groups (vide supra).

To confirm this, we have tried to reproduce the experimental spectrum by calculating different sets of crystal field parameters (CFP). Despite the high degeneracy of the Eu(III) ion (3003 levels), the 4f⁶ configuration is suitable to determine these CFP since the ⁷F multiplet is well isolated in the energy level scale: the ${}^{5}D_{0}$ to ${}^{7}F_{6}$ difference amounts to 12000 cm⁻¹, which implies that the ⁷F wave functions are relatively pure in ${}^{7}F_{J}$ states (>90%). Moreover, the crystal field operator operates only on states with same multiplicity; the ${}^{7}F_{J}$ multiplet is the only septet of the configuration, and therefore it is reasonable to perform calculations using the reduced base of the 49 $|^{7}F_{Jm}\rangle$ states.³¹ With use of Wybourne's formalism,³² the crystal field Hamiltonian can be written

$$H_{\rm c} = \sum_{kqi} B^k_{\ q} (C^k_{\ q})_i \tag{1}$$

The number of nonzero B^k_q CFP depends on the site symmetry. In our case, the low site symmetry, C_2 , along with the relatively few experimentally determined levels could lead to difficulties in fitting the spectrum. We have therefore tried to use the descending symmetry principle: D_{3h} was first considered as approximate symmetry in view of the above discussions, allowing us to determine four parameters; these were used as starting parameters to fit the data in the C_{2v} point group of symmetry, which is close to the actual C_2 symmetry.

In the D_{3h} symmetry, the expression for the crystal field Hamiltonian is

$$H_{\rm c} = B_0^2 C_0^2 + B_0^4 C_0^4 + B_0^6 C_0^6 + B_6^6 (C_6^6 + C_{-6}^6)$$
(2)

From the experimental splittings (Figure 5) it is easy to deduce a negative sign for B_0^2 as well as the magnitude of the B_0^4 parameter. The accurate determination of the four parameters is done by minimizing the root-mean-square (rms) deviation. The refinement leads to a final standard deviation of 18 cm⁻¹ (Table V). The sign of the B_6^6 parameter cannot be determined since opposite values do not modify the energies but only the labels of the irreducible representations; these latter cannot be obtained from measurements on powders

The crystal field potential is usually calculated considering the crystallographic axis as the z axis. However, it is possible to express this potential with respect to any other axis. In D_{3h} symmetry, one can therefore express the Hamiltonian (1) with respect to the C_2 axis, which corresponds to the crystallographic



Figure 6. Part of the excitation spectrum of $(Ph_4As)_2Eu(NO_3)_5$ at 77 K (analyzing wavelength 618 nm).



Figure 7. Part of the emission spectrum of (Ph₄As)₂Eu(NO₃)₅ at 77 K (excitation wavelength 396 nm).

axis. It will then be described by nine parameters corresponding to the CFP of the C_{2v} point group of symmetry: $H_1(C_2) = B^2 C_2^2 + B^2 (C_2^2 + C_2^2) + B^4 C_2^4 + C_2^4$

$$B_{2}^{\prime}(C_{2\nu}^{\prime}) = B_{-0}^{\prime}C_{-0}^{\prime} + B_{-2}^{\prime}(C_{-2}^{\prime} + C_{-2}^{\prime}) + B_{-0}^{\prime}C_{-0}^{\prime} + B_{-2}^{\prime}(C_{-2}^{\prime} + C_{-2}^{\prime}) + B_{-0}^{\prime}C_{-0}^{\prime} + B_{-0}^{\prime}(C_{-2}^{\prime} + C_{-2}^{\prime}) + B_{-0}^{\prime}C_{-0}^{\prime} + B_{-0}^{\prime}C_{-0}^{\prime}$$

These two potentials, $H_c(C_{2\nu})$ and $H_c(D_{3h})$, are identical, and the relationship between parameters of the same rank is described by a transformation matrix.³³ These parameters are used as starting parameters for the refinement in C_{2v} symmetry. We have determined two sets of parameters because of the uncertainty in the sign of B_6^6 (Table III). In both cases, the final rms deviations are small but the difference between them does not allow us to decide which is the sign of B_6^6 . In both cases too, the final parameters are not much different from those determined for a D_{3h} symmetry, which means that the adopted procedure of the descending symmetry is valid here. Moreover, the good agreement obtained with the D_{3h} parameters supports the assumption made above on the geometry of the coordination polyhedron, namely that it can be considered as arising from a distortion of a potentially more symmetrical arrangement. It should be stressed here that the more intense component of the ${}^{5}D_{0} \rightarrow {}^{7}F_{2}$ transition corresponds to the allowed transition $A_{1}' \rightarrow E$ in D_{3h} symmetry. The experimental and calculated levels of the ${}^{7}F_{j}$ manifold are compared in Table VI.

Emission Spectrum of $(Ph_4As)_2Eu(NO_3)_5$. The excitation and emission spectra were recorded at 77 K. They are partly

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Table VI. Comparison between Experimental and Calculated Sublevels for $(Ph_3EtP)_2Eu(NO_3)_s$ (cm⁻¹)

			calcd levels				
					20		
J	exptl levels	D,	h	sign	+a	sign	_a
0	0	0	A,'	0	A ₁	0	A
1	291	300	Α,'	294	B ₂	294	Β,
	422	428	Ε″	419	B ₁	425	B
	444	428	Ε''	442	A_2	437	A ₂
2		1023	E''	1016	A ₂	1009	Β,
	1046	1023	E''	1030	B ₂	1034	Α,
	1059	1069	A_1'	1071	A ₁	1063	A 1
	1095	1107	E'	1104	A ₁	1104	A ₁
	1134	1107	E	1123	Β ₁	1135	B ₂
3	1696	1700	A 2''	1702	A ₂	1701	B ₁
	1809	1829	E"	1820	В,	1813	B ₁
		1829	E E'	1831	A ₂	1834	A ₂
	1852	-1847	E F'	1858	R R	1860	
	1899	1877	Ă.'	1877	B ²	1883	R R
		2023	A_1''	2022	B_1^2	2022	A_2^2
4	2541	2516	A,''	2516	A ₂	2523	B ₁
	2601	2620	E'	2594	A	2593	A_1
	2629	2620	E'	2632	B ₂	2639	Α,
	2650	2635	A_1	2659	A_1	2655	B ₂
	2007	2699	E E''	2675	\mathbf{B}_{1}	2678	A ₂
	2093	2099		2720	A ₂ D	2/21	
	2995	3008	F'	2994	Δ	2985	R ^A 2
	3040	3008	Ē'	3042	B ₂	3047	A_1^2
5		3628	Α."	3623	A.	3618	В.
	3730	3741	E′	3736	В,	3737	A,
	3756	3741	\mathbf{E}'	3753	A,	3758	B ₂
	3785	3803	Е''	3791	A ₂	3782	A_2
	3819	3803	Ε''	3808	A ₂	3817	Β _ι
	3943	3959	A ₂	3952	B ₁	3956	B ₁
		3969	E'	3977	B ₂	3976	B ₂
		3707 1183	E A ''	398/ 1171	A ₂ P	398/	A ₂
	4183	4192	$\frac{A}{F'}$	4177		4181	R A
	4200	4192	Ē'	4202	B.	4205	$\frac{D}{A}^{2}$
			-		₽ ₂	1200	

^a Sign of the B_{6}^{6} parameter (see text).

displayed in Figures 6 and 7, respectively, and all the observed transitions are listed in Tables S3 and S4 (supplementary material) along with their assignments. Assignments for the emission spectrum were confirmed by selective excitations to the ${}^{5}D_{0}$ and to the ${}^{5}D_{1}$ levels, with either a dye laser (around 580 nm) or an argon laser (527 nm). The presence of only one sharp ${}^{5}D_{0} \rightarrow {}^{7}F_{0}$ transition (fwhh = 2.4 Å) is again consistent with one crystallographic site for the Eu(III) ion. The integrated and corrected relative emission intensities of the transitions from the ${}^{5}D_{0}$ level to the ${}^{7}F_{J}$ manifold are 0.4, 11.4, 100.0, 0.8, 9.5, 0.4, and 0.5% for J = 0-6, respectively. This intensity pattern is quite similar to that found for the previous compound, the main differences lying in the intensities of the transitions to the J = 0 and J = 5 levels, which are much larger in the tetraphenylarsonium complex. A scheme of the crystal field splittings for the ${}^{5}D_{J}$ and ${}^{7}F_{J}$ manifolds is presented in Figure 8. Contrary to what was observed for the triphenylethylphosphonium complex, the pattern of these splittings along with the presence of a moderately intense 0-0 transition does not reflect a potentially higher symmetry for the europium site than the crystallographic C_2 symmetry. We have therefore calculated the crystal field parameters in C_{2v} symmetry using the sets determined for $(Phe_3EtP)_2Eu(NO_3)_5$ as starting parameters.³⁴ A difficulty arose because the 5D_0

Table VII. Calculated Crystal Field Parameters (CFP, cm^{-1}) for (Ph₄As)₂Eu(NO₃)₅ in $C_{2\nu}$ Symmetry

· •	· 4	20 .	•		
СГР	sign + ^a	sign – ^a	CFP	sign + ^a	sign – ^a
$ \begin{array}{r} B^{2}{}_{0} \\ B^{2}{}_{2} \\ B^{4}{}_{0} \\ B^{4}{}_{2} \\ B^{4}{}_{4} \\ B^{4}{}_{4} \end{array} $	$232 \pm 9 \\ -26 \pm 8 \\ -352 \pm 19 \\ -135 \pm 18 \\ -572 \pm 17$	326 ± 9 34 ± 5 -120 ± 19 -345 ± 11 -469 ± 13	$\begin{array}{c}B^{6}{}_{0}\\B^{6}{}_{2}\\B^{6}{}_{4}\\B^{6}{}_{6}\\\sigma\end{array}$	$\begin{array}{r} -2223 \pm 24 \\ 973 \pm 21 \\ -750 \pm 18 \\ -437 \pm 20 \\ 15.4 \end{array}$	$1481 \pm 19 \\ -1560 \pm 13 \\ +225 \pm 16 \\ -696 \pm 15 \\ 12.0$

^a Refers to the two starting parameter sets; cf. Table V.

Table VIII. Comparison between Experimental and Calculated Sublevels for $(Phe_4As)_2Eu(NO_3)_5$ (cm⁻¹)

			vels, C _{2v}			
J	exptl levels	sign	+ a	sign -	_a	
0	0	0	A ₁	0	A ₁	
1	300	319	Β,	299	В,	
	362	347	B	351	B ₁	
	429	425	A ₂	440	A ₂	
2		93 0	A ₂	931	A ₂	
	1007	1000	B	993	B ₁	
	1026	1022	B ₂	1019	A 1	
	1068	1064	A ₁	1085	A ₁	
2	1082	1102	A,	1090	В ₂	
3	1670	1675	A ₂	1672	B ₁	
	1794	1789	A ₂ B	1794	B ₁	
	1803	1819	\mathbf{A}_{1}	1806	B.	
	1836	1821	В,	1841	B,	
	1863	1855	B ₂	1843	A ₁	
	1947	1949	B	1955	A ₂	
4	2535	2513	A ₂	2514	B ₁	
	2568	2557	A ₁	2552	A ₁	
	2582	2575	\mathbf{A}_{1}	2589	A ₁	
	2601	2600	B ₂	2599	B ₂	
	2636	2635	B ₁	2643	B ₁	
	2033	2082	A ₂ B	2080	A_2	
	2932	2950	A_{1}	2936	B.	
	2994	2975	B,	2981	\tilde{A}_{1}^{2}	
5		3706	A.	3695	В.	
	3767	3790	B ₂	3765	\mathbf{A}_{1}^{-1}	
	3829	3819	A ₁	3833	B	
	3845	3846	A ₂	3844	B ₁	
	3872	3874	B	3861	A ₂	
		4005	A ₂ B	3980	A ₂ B	
		4033	B.	4040	B.	
		4281	B	4258	Å,	
		4286	A ₁	4262	B ₂	
		4287	B ₂	4273	A_1	
6	4929	4959	B ₂	4952	A ₁	
		4968	A ₂	4954	B ₁	
	5018	5014	B	5017	A ₂	
	5077	5015		5026	В ₂	
	5096	5083	A ₁ B	5075		
	0070	5086	\mathbf{A}_{2}^{1}	5106	B.	
	5117	5131	A ₁	5109	Ă,	
	5127	5139	A ₂	5116	B ₂	
	5142	5140	B ₂	5153	\mathbf{A}_{1}	
	5181	5157	B ₁	5174	B ₂	
	5215	5218	В ₂	5201	A 2	
	5215	5219	A_1	3211	A	

^a Sign of the B_{6}^{6} parameter (see text).

→ ${}^{7}F_{2}$ transition, when measured with high resolution under laser excitation, displayed too many components as already observed for other complexes.^{18,35} We have considered the two closely spaced multiplets at 616.37 and 616.59 nm and at 617.10, 617.20, and 617.27 nm as representing only one

⁽³⁴⁾ A calculation in D_{3h} symmetry results in a poor fit, with $\sigma = 27$ cm⁻¹; the parameters are very similar to those obtained for the triphenyl-ethylphosphonium complex.

⁽³⁵⁾ Judd, B. R. Phys. Scr. 1979, 21, 526.



Figure 8. Crystal field splittings of $(Ph_4As)_2Eu(NO_3)_5$, in cm⁻¹, from excitation and emission spectra at 77 K. The vertical scale for 7F_4 has been reduced by a factor of 2. The numbers between parentheses are the barycenters of the levels calculated from the observed transitions.

transition. The exact origin of these "additional" splittings is presently under investigation in our laboratory. The results of the spectrum simulation are presented in Tables VII and VIII. Compared to the previous calculation, the overall fit is not as good, although it remains quite acceptable in view of the actual lower symmetry C_2 .³⁶ The set of parameters with a positive B_0^6 value gives a final standard deviation significantly lower than the other set, and it also better reproduces the splittings of the J = 1 and J = 2 levels. The corresponding set for the tetraphenylarsonium complex also has a smaller σ and can therefore be considered as more adequate. In general the crystal field parameters for the two pentanitrates have quite similar values, especially the three larger ones, B_{0}^6 , B_{2}^6 , and B_{6}^6 . The only exceptions are B_{2}^2 and B_{4}^6 , which vary by a factor of 8–10 but which are particularly small for Phe₄AsP⁺.

Emission Spectrum of (Me_4N)_2Eu(NO_3)_5. Depending on the crystallization and drying procedures used, one, two, or even three ${}^5D_0 \rightarrow {}^7F_0$ transitions are observed for this compound, at 580.2, 579.8, and 579.1 nm, respectively. Perfectly anhydrous samples always display more than one 0–0 transition, while hydrated samples have only one. The crystal field splittings of these latter exhibit a pattern reflecting a higher pseudosymmetry, namely C_{3h} or D_{3h} . We have, however, not analyzed it in view of the uncertainty on the structure of the hydrated compound.

Fluorescence Lifetimes. The fluorescence lifetimes are equal to 0.74 (2) and 0.67 (1) ms for the triphenylethylphosphonium and the tetraphenylarsonium compounds, respectively. They

are temperature independent between 77 and 300 K, which means that the nonradiative processes involve a large number of vibrational quanta.³⁷

Conclusion

The detailed analysis of the emission spectra and of the structural data of pentakis(nitrato)europate(III) ions performed in this paper allows us to draw some conclusions on the coordination polyhedron around the Eu(III) ion. The bidentate nitrates induce a low symmetry, C_2 ; however, the nitrogen atoms of these ligands define a more symmetrical arrangement, namely an almost undistorted trigonal bipyramid. For the triphenylethyl compound, the symmetry of this arrangement is reflected in some of the crystal field splittings and the emission spectrum of the Eu(III) ion is reasonably well simulated by a set of four crystal field parameters defining an Hamiltonian in D_{3h} symmetry. This means that in these essentially ionic compounds the Eu(III) ion tends to realize the more compact and more symmetrical arrangement of its ligands. Moreover, the analysis of the crystal field splittings may be used to assess the extent of the distortion of a coordination polyhedron, as indicated by the poorer fit obtained with the tetraphenylarsonium compound, the crystal field splittings of which do not exhibit a pattern reflecting a higher pseudosymmetry. The relative intensity of the forbidden ${}^{5}D_{0}$ \rightarrow ⁷F₀ transition seems to confirm the conclusions drawn from the crystal field splitting analysis: this transition has an extremely low intensity in (Ph₃EtP)₂Eu(NO₃)₅, which has the less distorted arrangement, whereas it has appreciably more intensity in the spectra of $(Ph_4As)_2Eu(NO_3)_5$. On the other hand, the extreme sensitivity of the emission spectra to the environment of the Eu(III) ion makes it difficult to recognize the entity $[Eu(NO_3)_5]^{2-}$ solely on the basis of its emission spectrum.

The analysis performed here provides good starting crystal field parameters for a detailed simulation of the emission spectra via the method of descending symmetry. This may be of interest in quantifying the optical properties of potential phosphors for which only powdered samples are available.

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Supplementary Material Available: Excitation spectral data (Tables S1 and S3) and emission spectral data (Tables S2 and S4) (7 pages). Ordering information is given on any current masthead page.

⁽³⁶⁾ One should keep in mind that these CFP were calculated by assuming structural and spectroscopic similarities between the two [Eu(NO₃)₅]²⁻ moieties; it is likely that other CFP sets can be found, which would eventually also yield a good fit of the spectrum.

⁽³⁷⁾ Auzel, F. Phys. Rev. B: Solid State 1976, 13, 2809.